

Closures for Meso-scale Models of Dense
Suspensions

The Euler-Lagrange Perspective

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with contributions from

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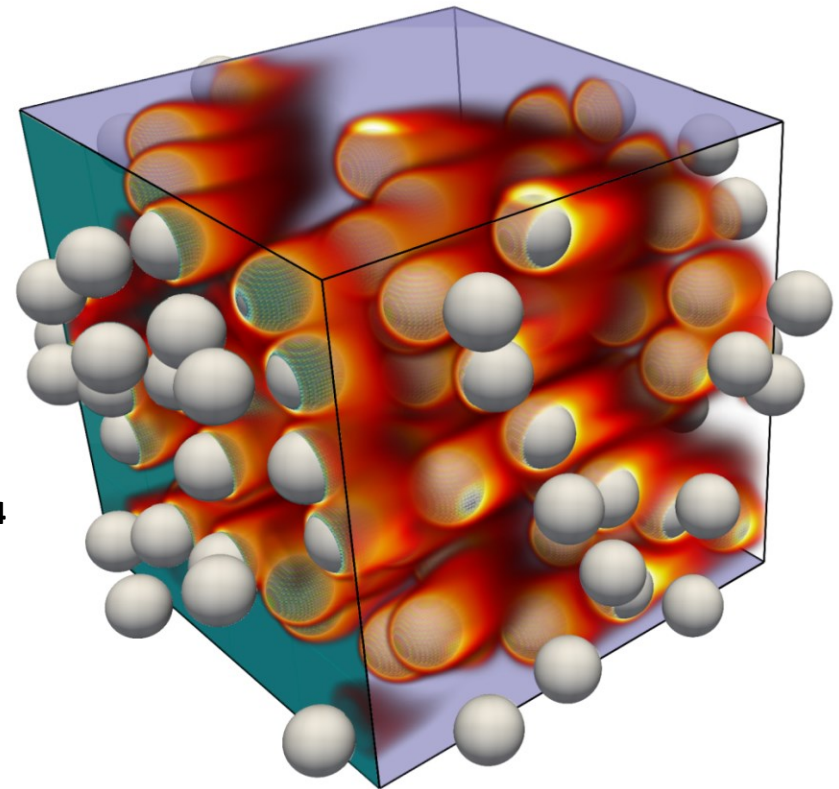
A. Ozel,² J. Kohlemainen,² S. Sundaresan,²

C. Goniva,³ A. Singhal,⁴ S. Cloete,⁴ S. Amini⁴

²Princeton University, N.J.

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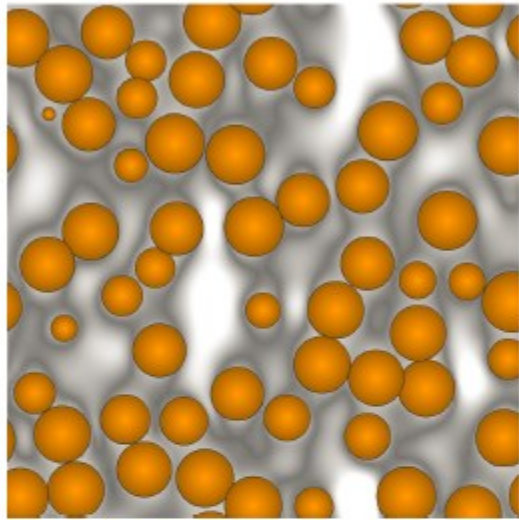
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Example Application of CxD

I – MICRO

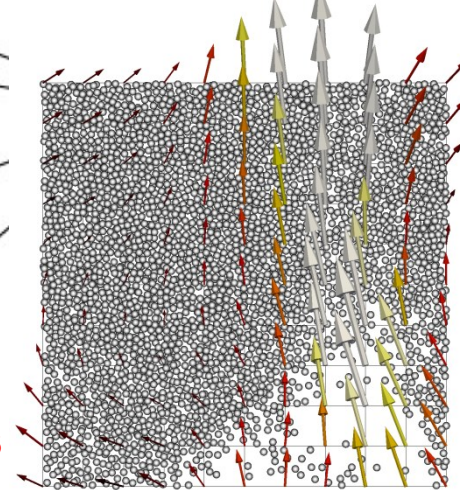
$\sim 50\mu\text{m}-\text{mm}$



II – MESO

$\sim \text{mm}-\text{cm}$

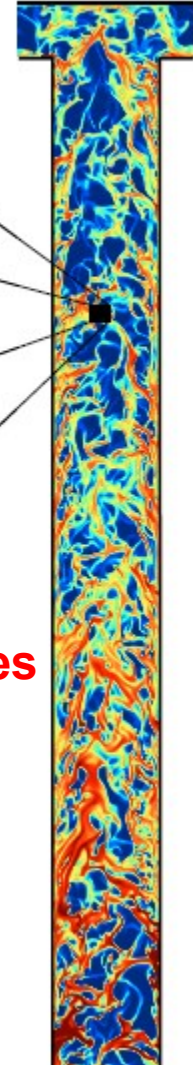
closures



closures



$\sim \text{cm}-\text{m}$



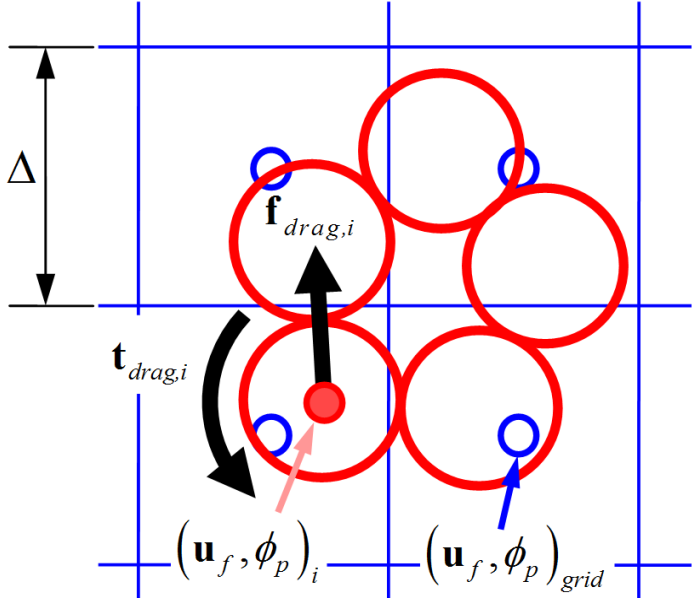
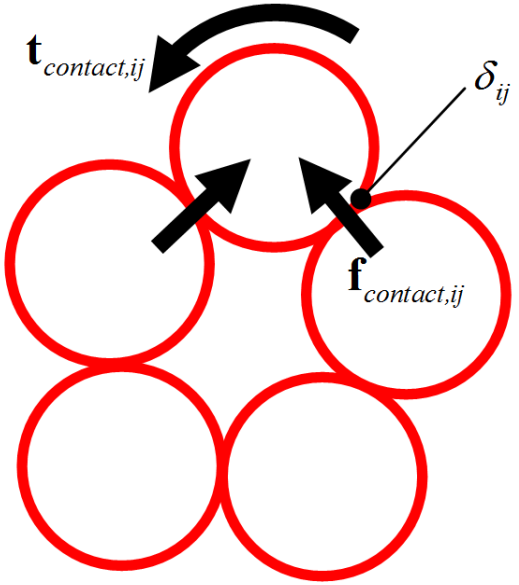
III MACRO

- **Meaningful** reaction kinetics must be fed into “**micro-scale**” models
- **Parameter Screening** at the micro scale to answer “*what matters*” (i.e., pore-scale diffusion, reaction, etc.)
- **Continue with meso and macro scale** if necessary.
- **Need closures!**

Closures at the Meso Level

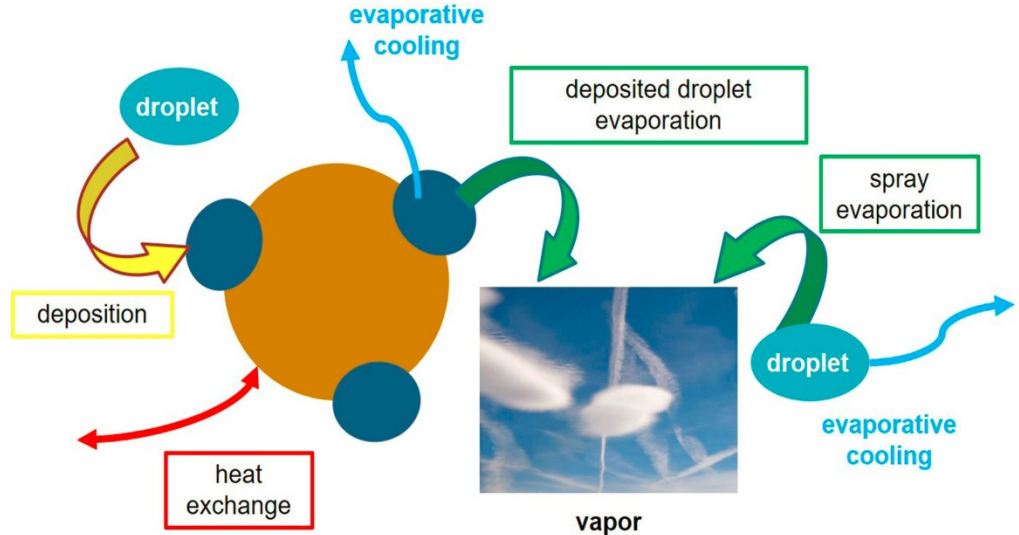
Flow

- Contact+cohesive forces and torques **per contact**
- Fluid-Particle interaction (drag) forces and torques **per particle**



Scalar Transport

- Heat and mass transfer rates (Nusselt/Sherwood numbers) **per particle**
- Dispersion rates **(fluid phase)**
- Filtration rates **per particle**
- Liquid transfer rates **per contact**

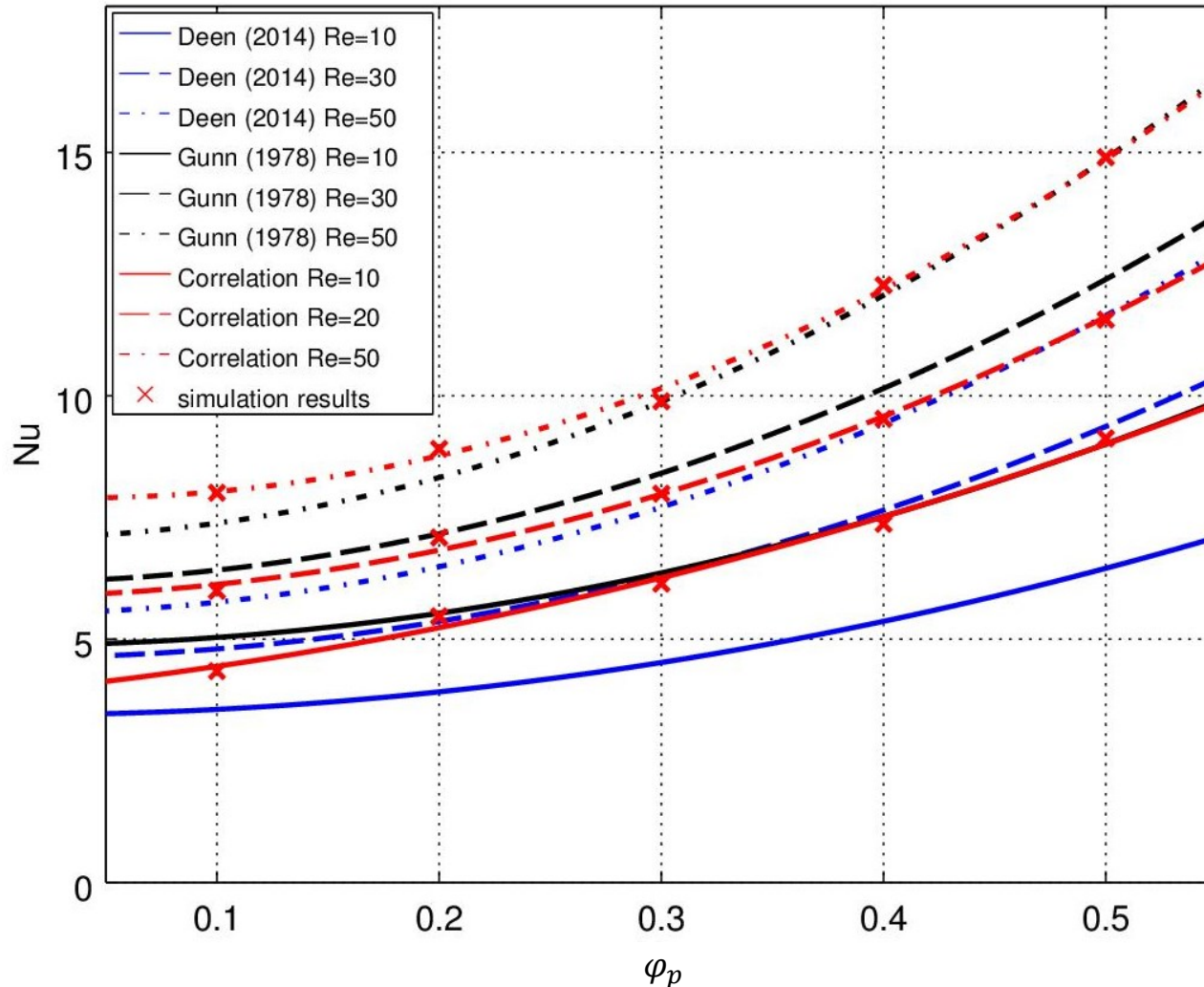


[2] M. Askarhishahi et al., *AIChE J* (2017) 63:2569-2587

- Part I **The Bad**
(...things done *incorrectly* in the past)
- Part II **The Hope** (...present research)
- Part III **The Future** (...most likely ‘The Good’)

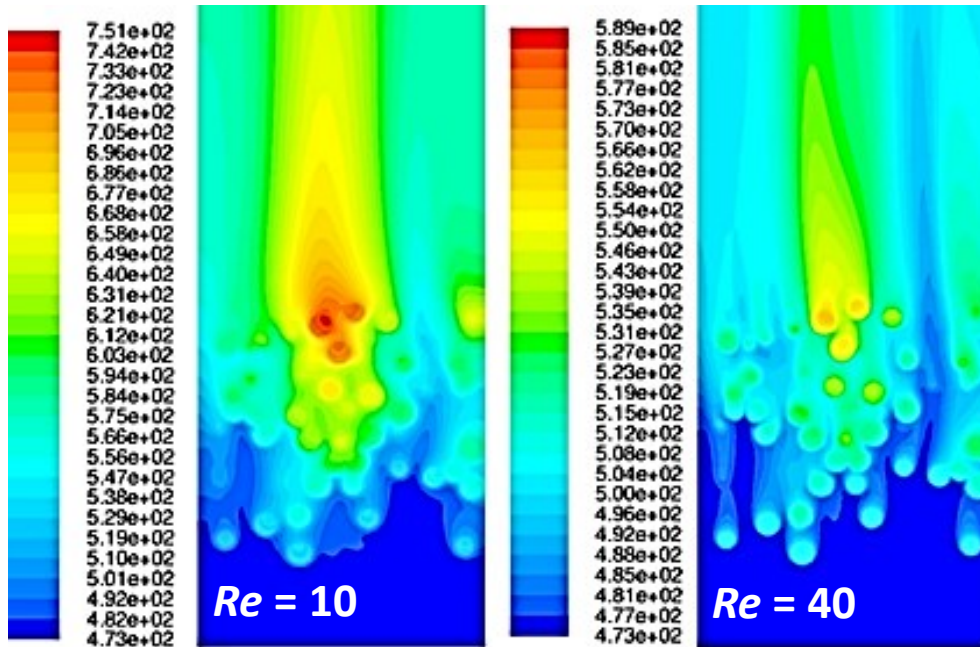
The Bad

The Correlations



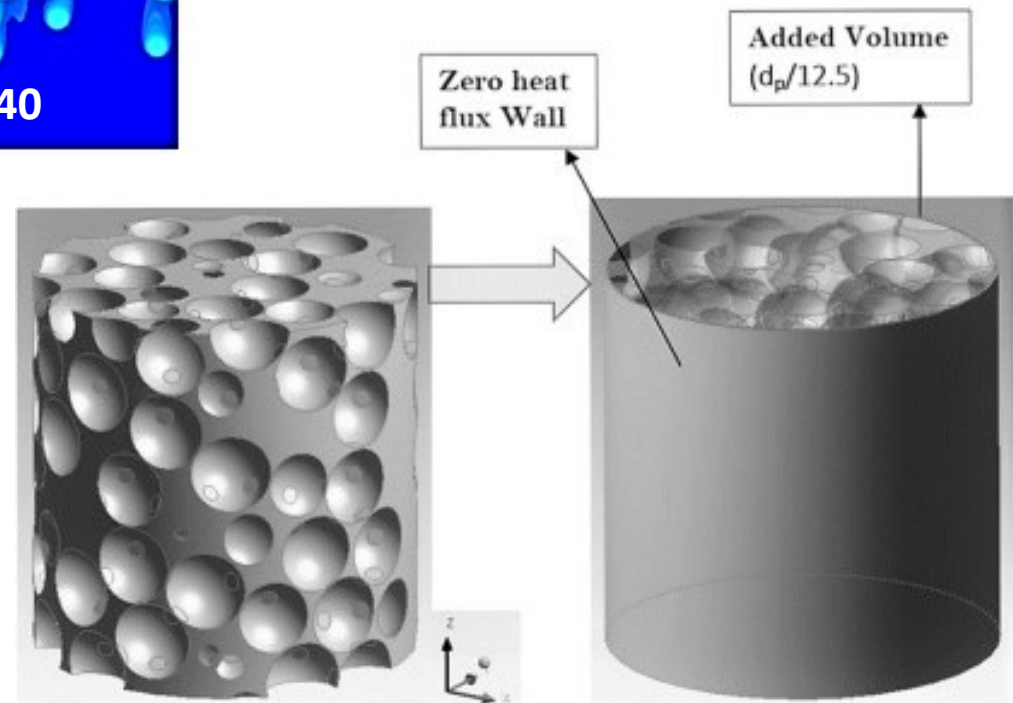
- Many exist for the **mean (i.e., an average over many particles)**
- The **fluid's cup-mixing** and **local mean temperature** are confused. Cup-mixing temp.: **okey for bed-average Nusselt numbers**, but in Euler-Lagrange models this quantity is **NOT known [3]!**
- **Correlations** are often **“over fitted”** in regimes where this is unnecessary (e.g., low Re, high ϕ_p)

The Cylinder



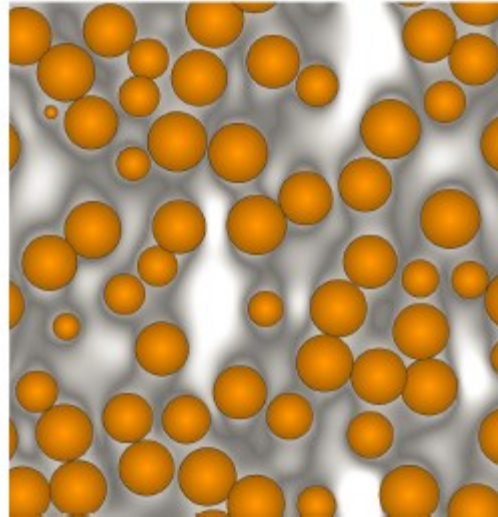
- Regions can be “cut out”: this is **cumbersome (meshing!)**
- **Wall distance and wall curvature effects** are mixed up

- Correlations are valid for the **mean** and far away from **walls**
- **Confidence intervals** for parameters are not provided
- Computational domains are often **too small**

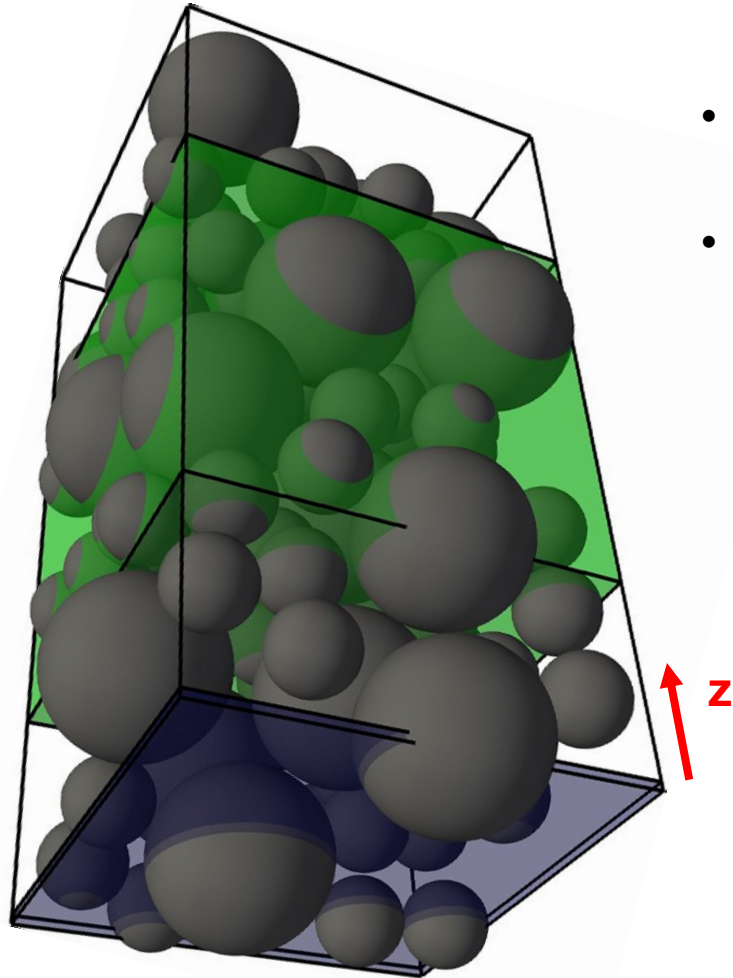


The Hope

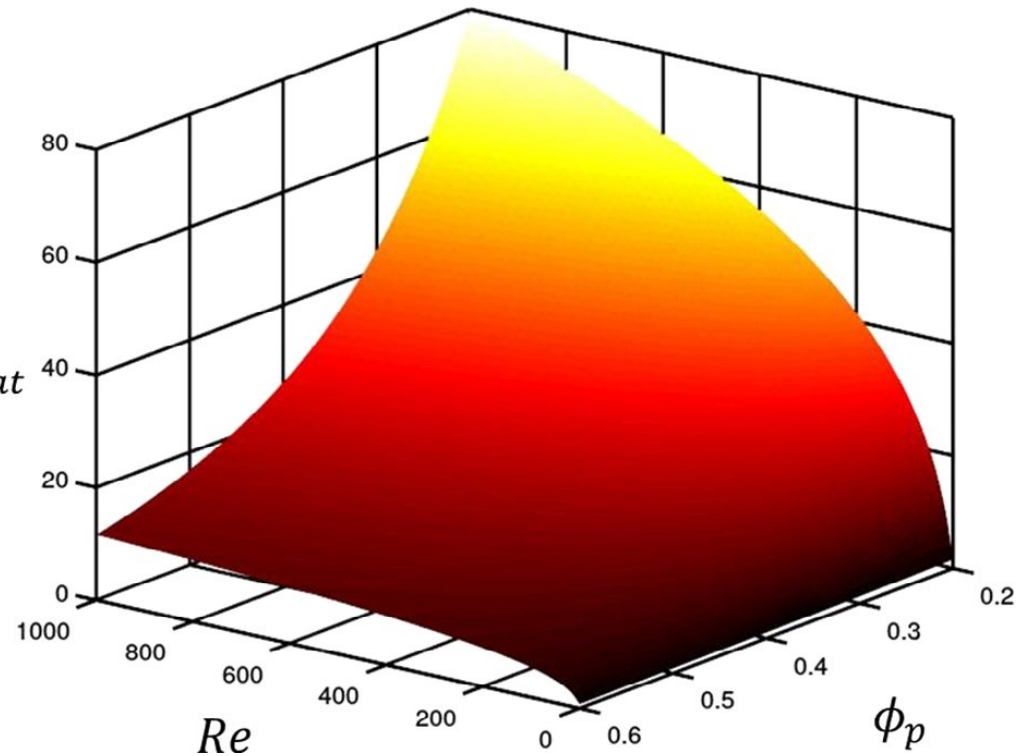
Ila – Towards Improved Closures



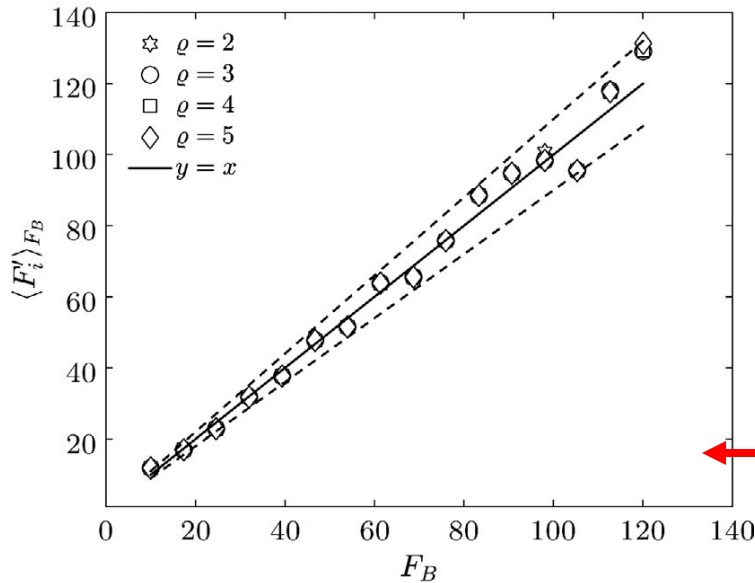
Saturation



- For small Re and high $\phi_p \rightarrow$ fluid phase is **quickly saturated** with the transferred quantity (**i.e., small z_{sat}**)
- Fluid **field quickly relaxes** to equilibrium value provided at particle surface
- In a meso-scale simulation, **Nu would NOT matter!**

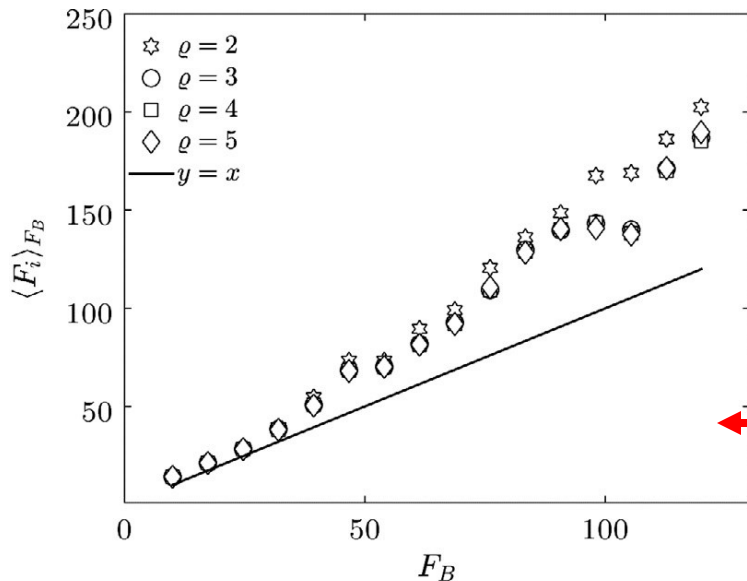


Bi-Disperse Systems: Drag Coefficient



- One **cannot simply re-scale** the fluid-particle interaction force (with $1 - \phi_p$) to extract the drag force in bi- (and poly) disperse suspensions
- Fortunately, this can be **“repaired”**

Municchi and Radl (**simple re-scaling**) versus Beetstra et al. (**simple re-scaling**)

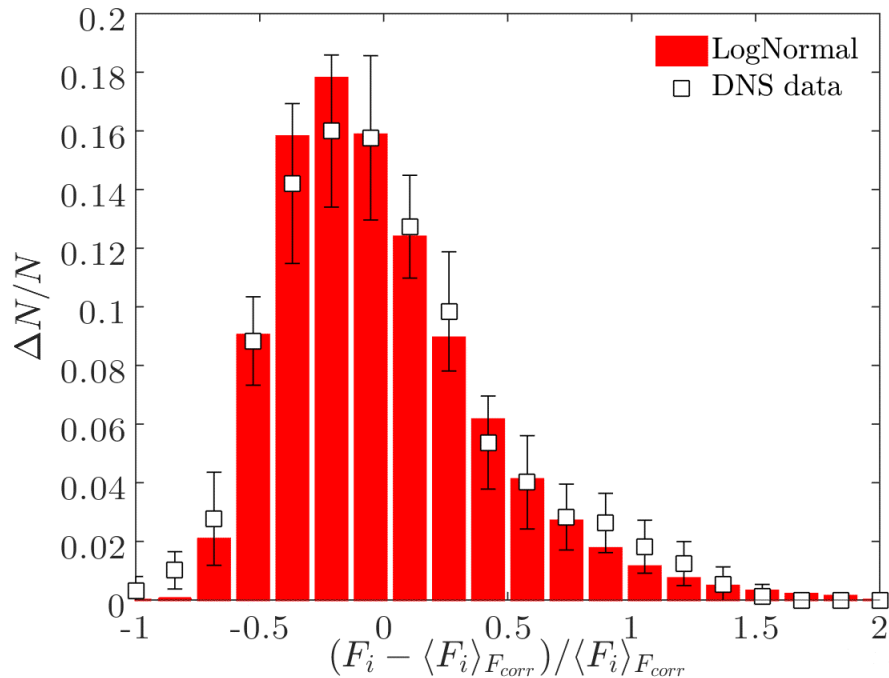


$$\mathbf{f}_{drag,i} \equiv \mathbf{f}_i - \mathbf{f}_i^{\nabla p^e}$$

- \mathbf{f}_i : Total force acting on particle i
- \mathbf{f}_i^d : Drag force acting on particle i
- $\mathbf{f}_i^{\nabla p^e}$: Force due to mean pressure gradient

Municchi and Radl (**correct pressure gradient handling**) versus Beetstra et al. (**simple re-scaling**)

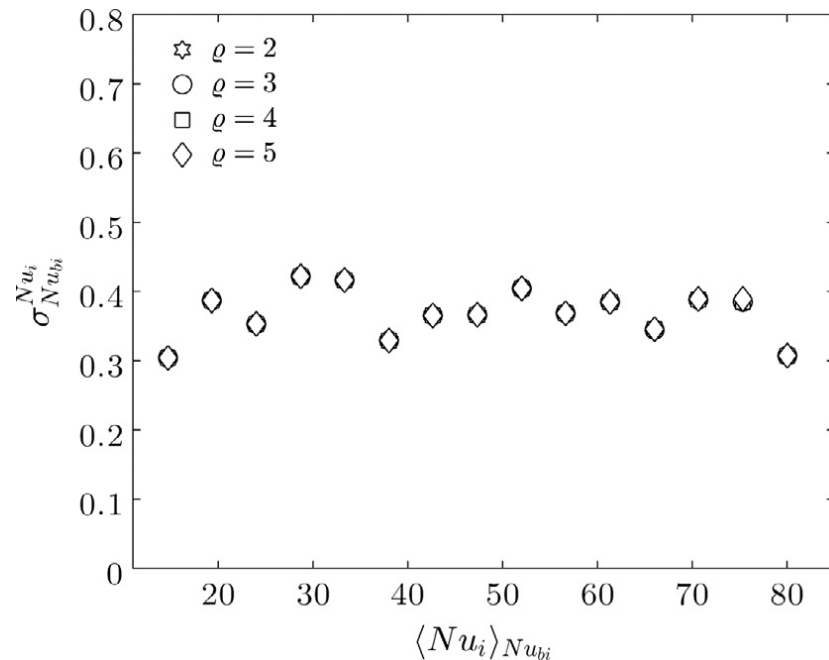
Drag Coefficient



- Previous work [6] on per-particle drag variation attempted to model the **total fluid-particle force** (with moderate success)
- However, when using a correctly-defined drag coefficient: the **scaled variance for the drag coefficient** is approximately constant: **simple closure possible!**
- Particle-individual deviations can be approximated using a **Log-Normal distribution**

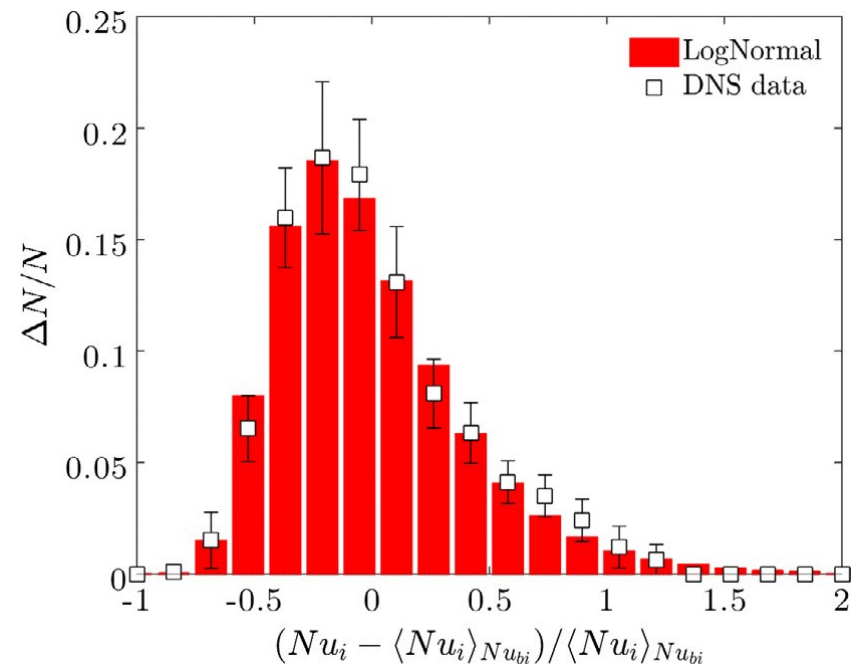
Bi-Disperse Systems: Mean versus Per-Particle

Nusselt Number

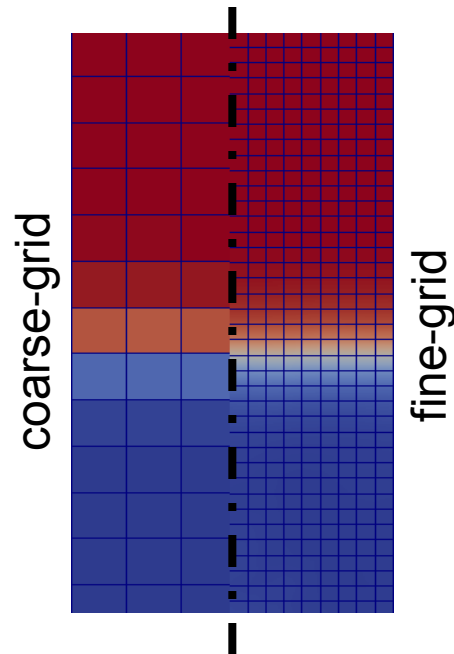


- Particle-individual deviations again follow a **Log-Normal distribution**, which is a bit more peaked.

- Same as for the drag coefficient: **scaled variance for the Nusselt number is approximately constant: simple closure possible!**



Ila – Towards Improved Voidage Reconstruction



Eulerian versus Lagrangian

Base case: fine grid

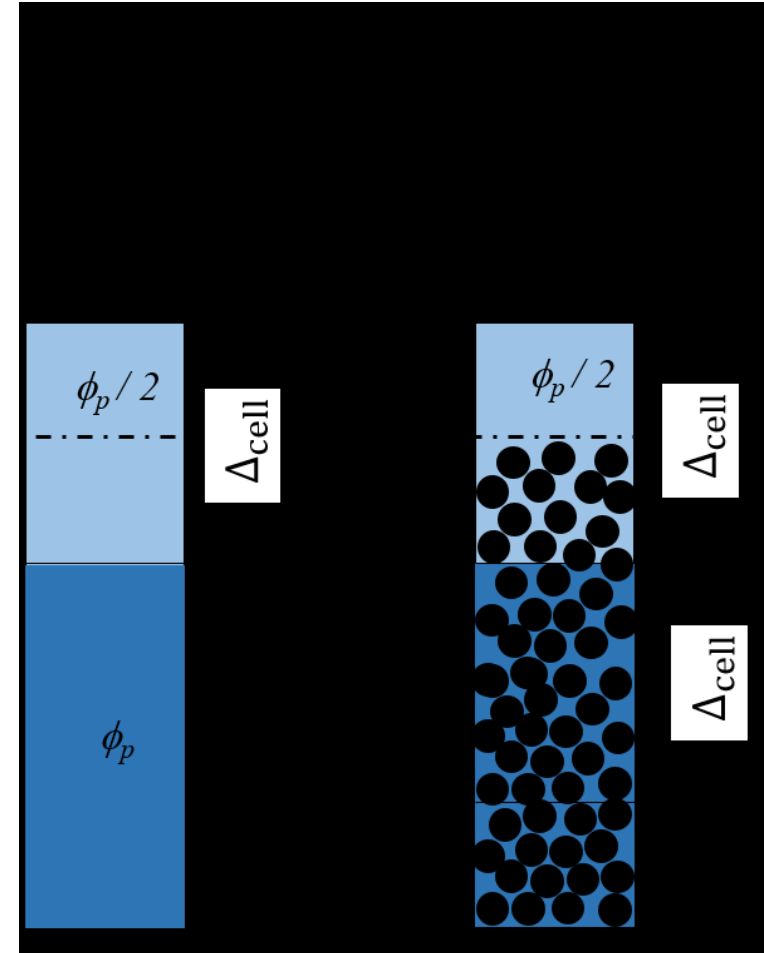
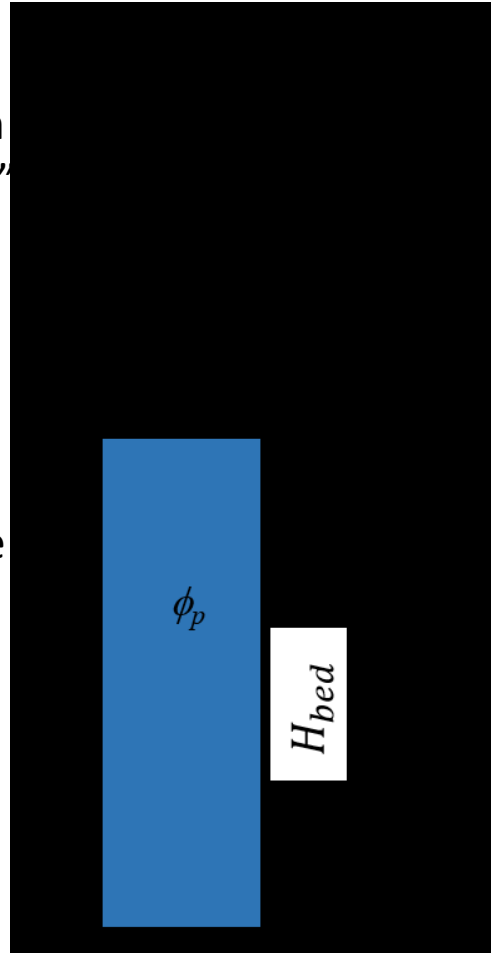
grid aligned with the jump in the voidage profile (“perfect” solution)

Case 1: “coarse Eulerian grid”

the jump in voidage profile at the centre of the interface cell

Case 2: “coarse Lagrangian grid”

the voidage is linearly interpolated at each particle position

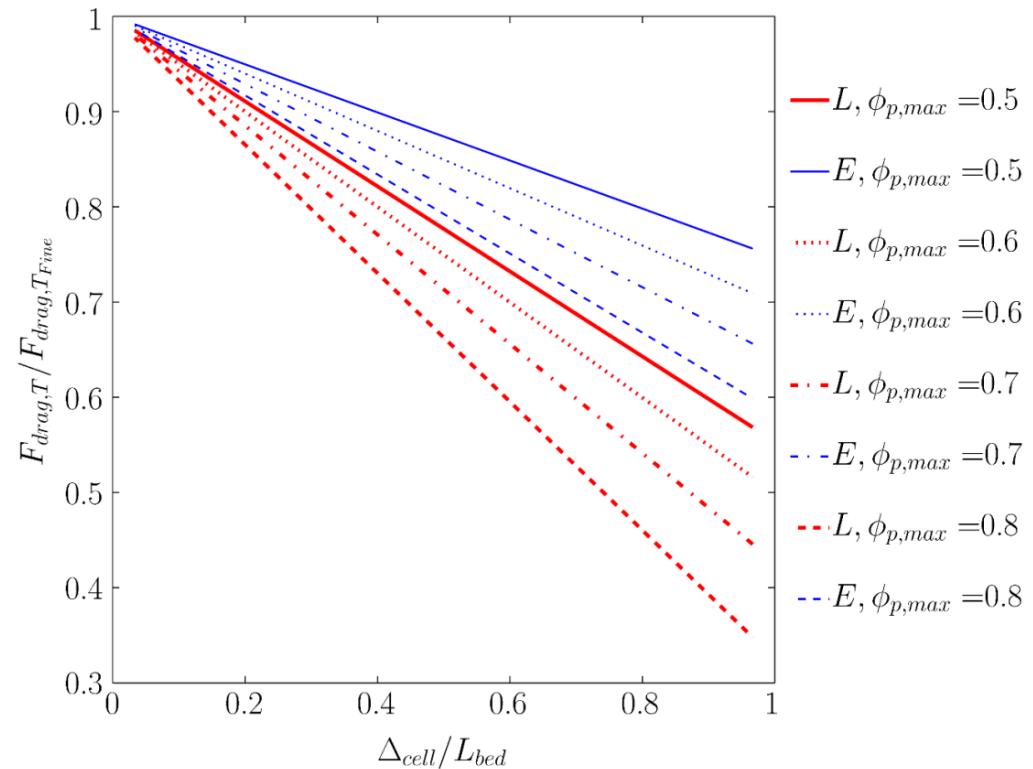


Eulerian versus Lagrangian

Local voidage is
*overpredicted on
average*

*Underprediction of
exchange coefficients*

*Voidage correction is
needed!*



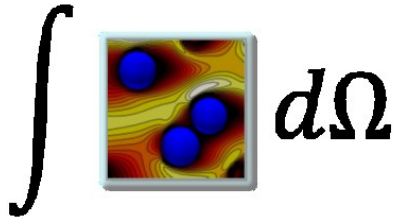
Lagrangian approach results in a **larger error** compared to Eulerian approach!

The Future

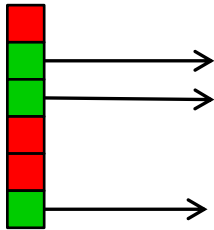
Coherent Toolsets

Post-Processing Utilities (e.g., CPPPO)

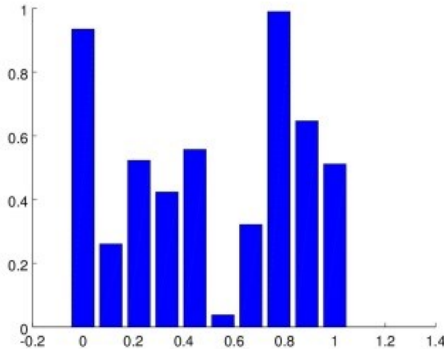
A Typical Set of Operations



Filtering of fluid and particle data, including **variance calculation**



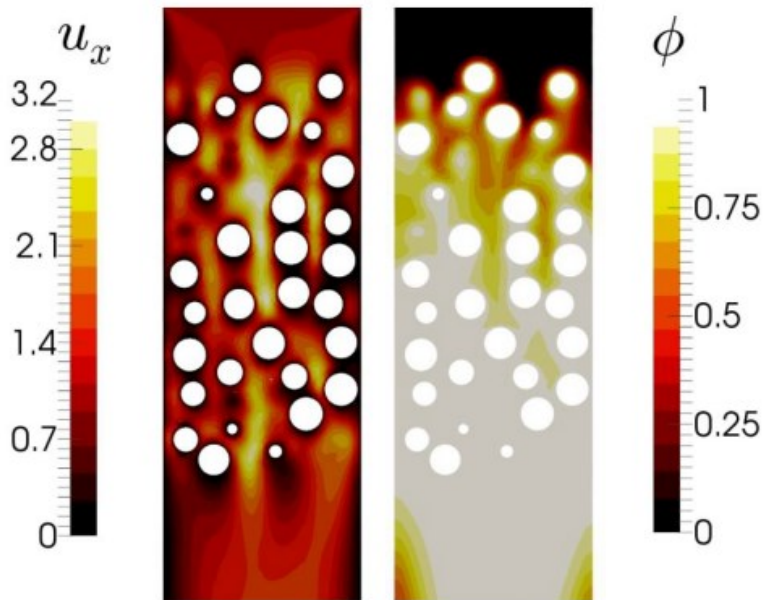
Sampling of filtered data (**defined at runtime**) and their derivatives with statistical **biasing** (e.g., limiters)



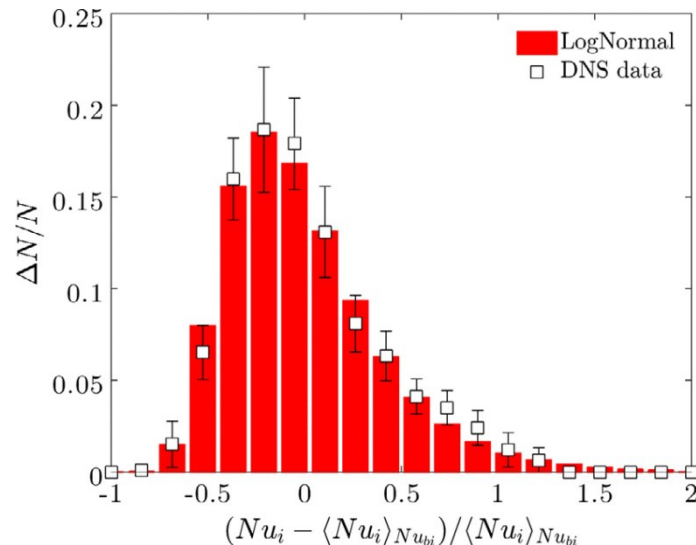
Binning of sampled data using running **statistics**

Coherent Toolsets

Post-Processing Utilities (e.g., CPPPO)

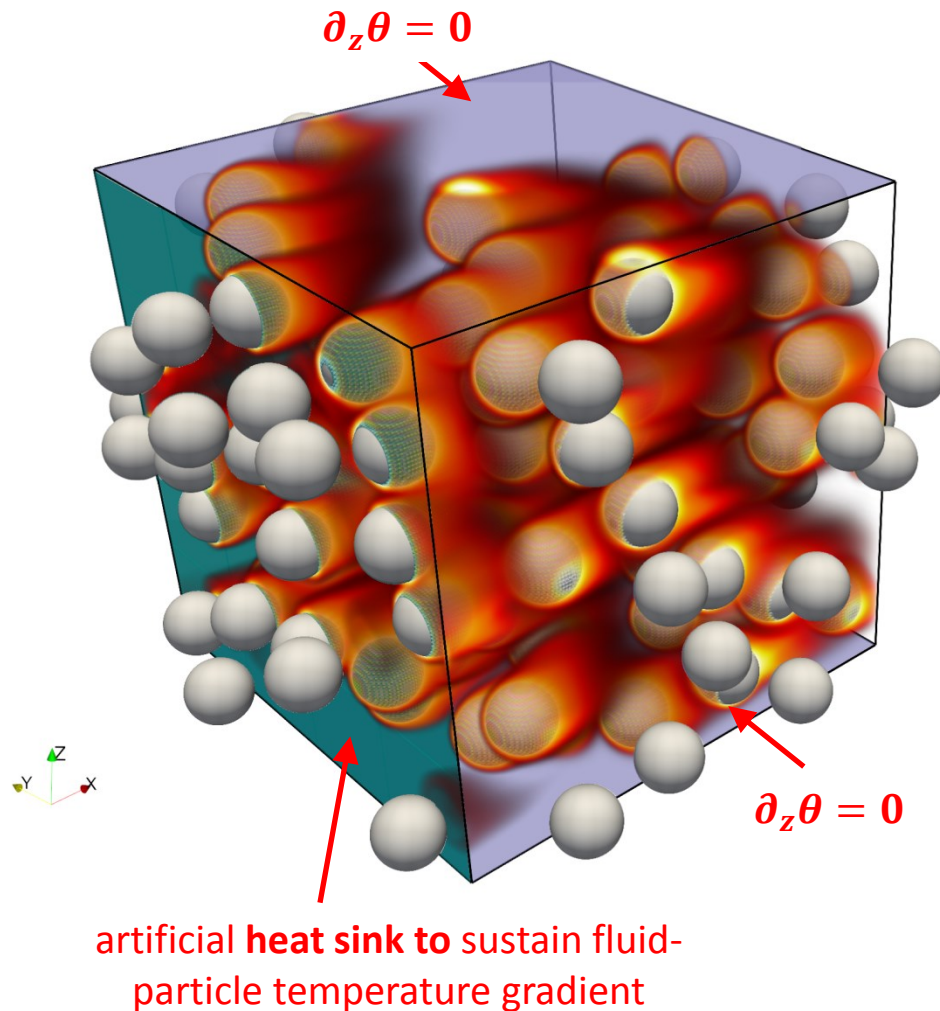


- **Support theory** (NOT mindless parameter fitting!): test hypothesis, supply data to establish closure, etc..
- Faster evaluation of filtered quantities desirable (**differential filtering**).
- Exploration of a wider array of raw data sources (“**embedded** DNS boxes”, “**forcing**”) desirable → **database of filtered statistics**



Particle-Resolved DNS to identify Modeling Needs

Boundary conditions: temperature field



- Particle bed generated via bi-axial **compaction** in the xy plane using LIGGGHTS®
- Flow and temperature fields are solved in a xy **periodic domain**. **Particles are isothermal.**
- **CFDEM®Coupling** to solve the governing equations for the continuum phase
- Particles are represented by forcing terms in the governing equations, **Hybrid Fictitious Domain-Immersed Boundary** method

Walls

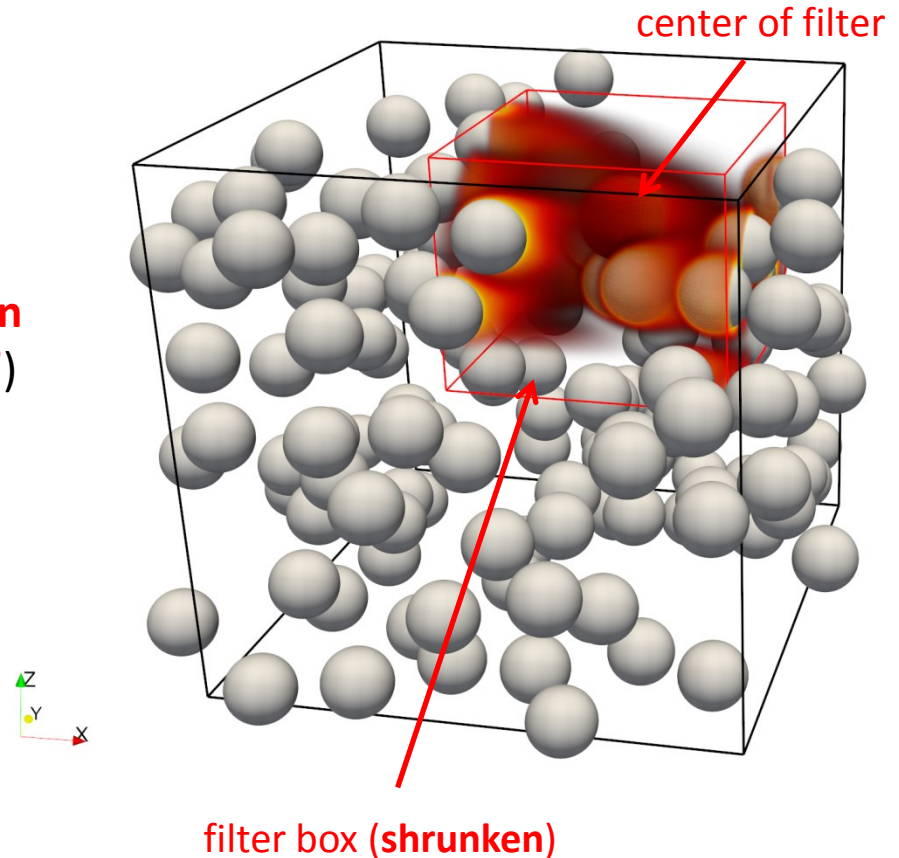
Particle-Resolved DNS to identify Modeling Needs

- We make use of the filtering toolbox **CPPPO** to spatially average (“filter”) the continuum phase properties around **each particle**

$$\varrho = \frac{L_{filter}}{d_p} \quad \begin{array}{l} \text{Dimensionless} \\ \text{filter size} \end{array}$$

- CPPPO is also employed to draw more “conventional” statistics (e.g., **profiles in wall-normal direction, “pancake filter”**)
- Filter boxes are **shrunk in the vicinity of wall boundaries**, same as done for wall bounded single phase turbulent flow

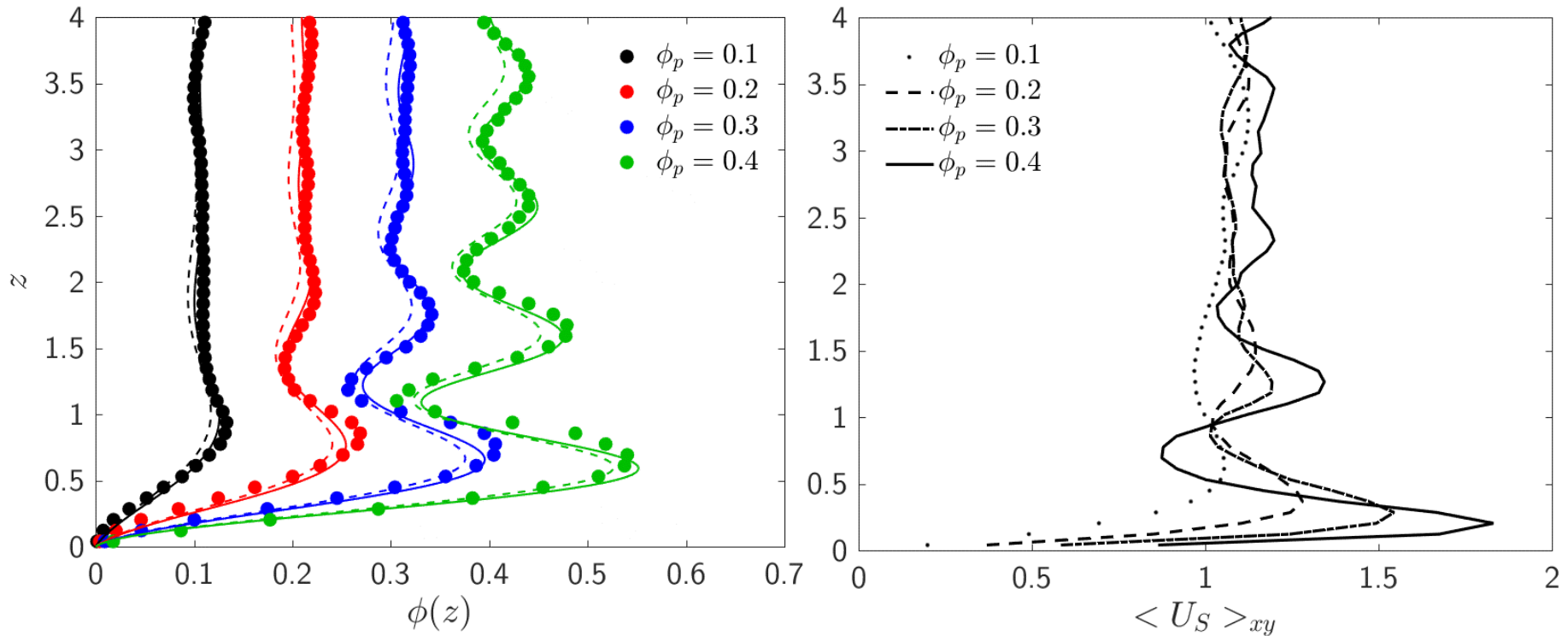
Lagrangian filtering: wall particles



Walls

Particle-Resolved DNS to identify Modeling Needs

Local Voidage and Speed

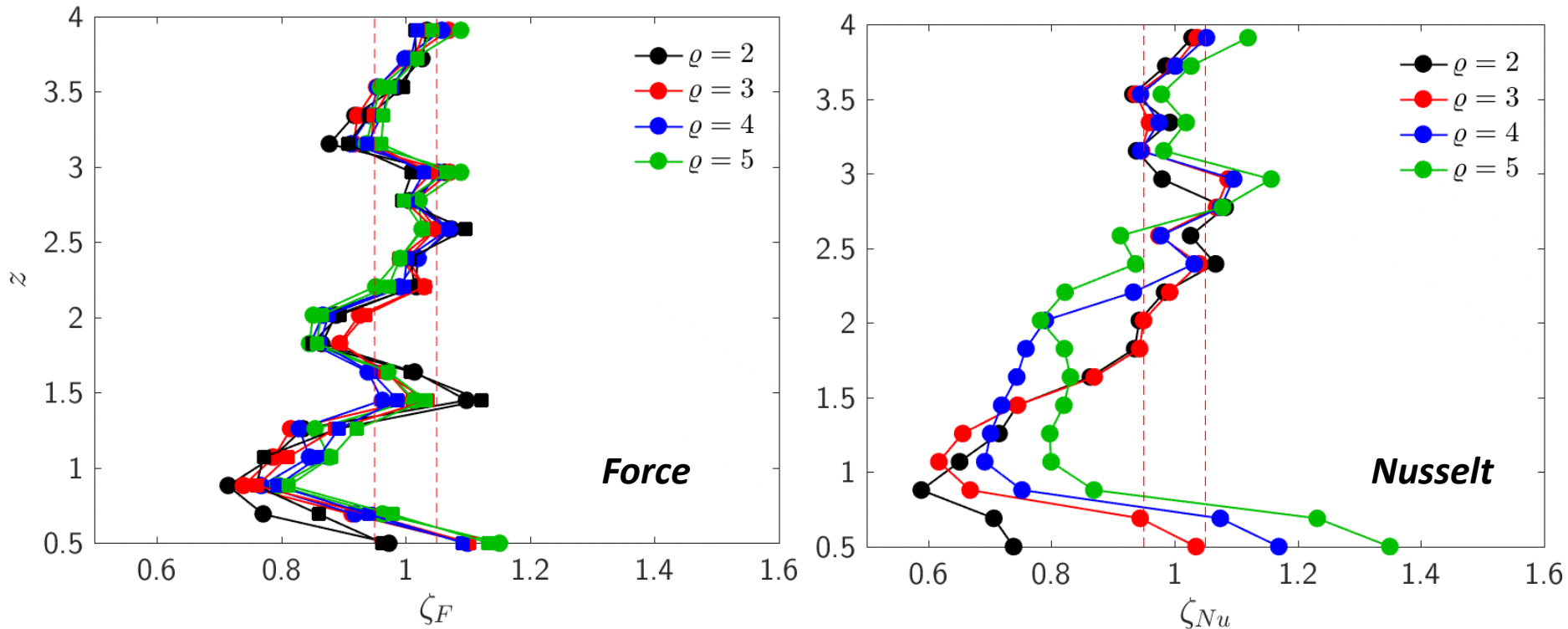


- General **correlation proposed for $\phi(z)$**
- Fluid speed **fluctuates strongly**, but with **small wavelength** \rightarrow we expect a **filter-size independent near-wall correction**

Walls

Particle-Resolved DNS to identify Modeling Needs

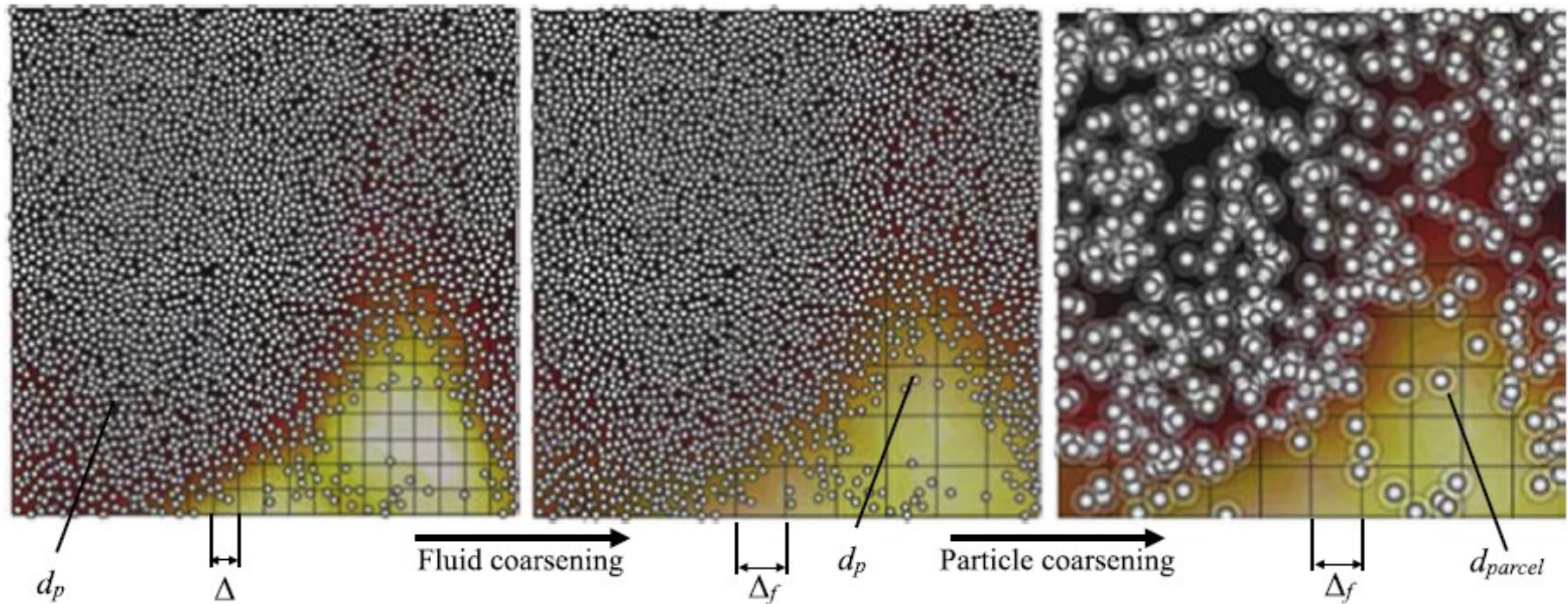
Local Drag Correction and Nusselt Number



- $\langle \phi_p \rangle = 0.4$: substantial **negative drag correction** for “2nd layer” particles
- For the Nusselt number, the situation is **more complex (due to temperature profile!)**, and even higher **(mixed) heat flux corrections** are necessary

Faster Simulations

Deriving Closures for Large(r) Scale-Models

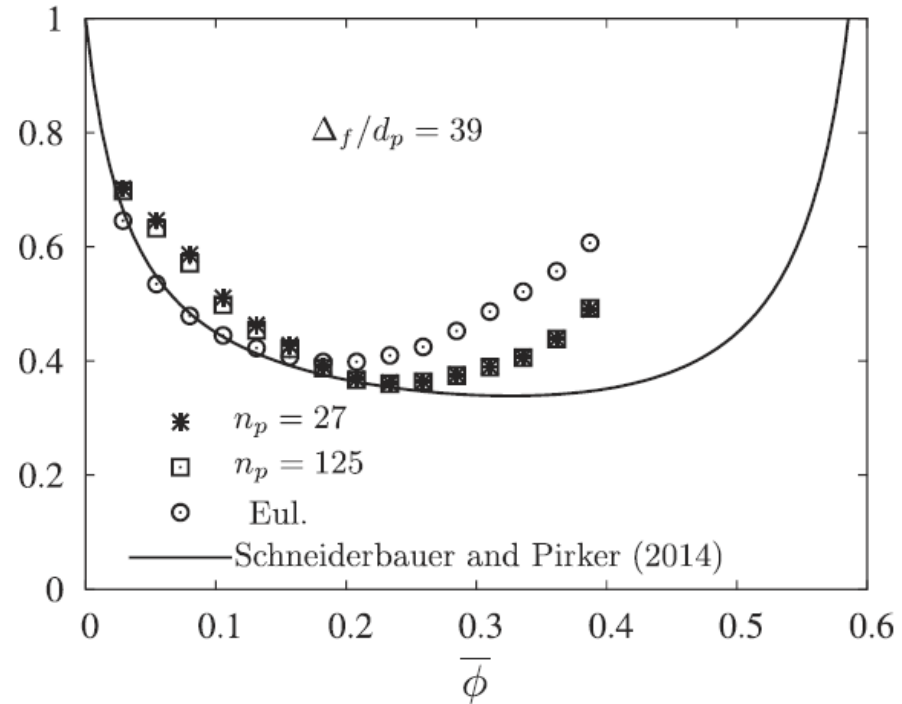
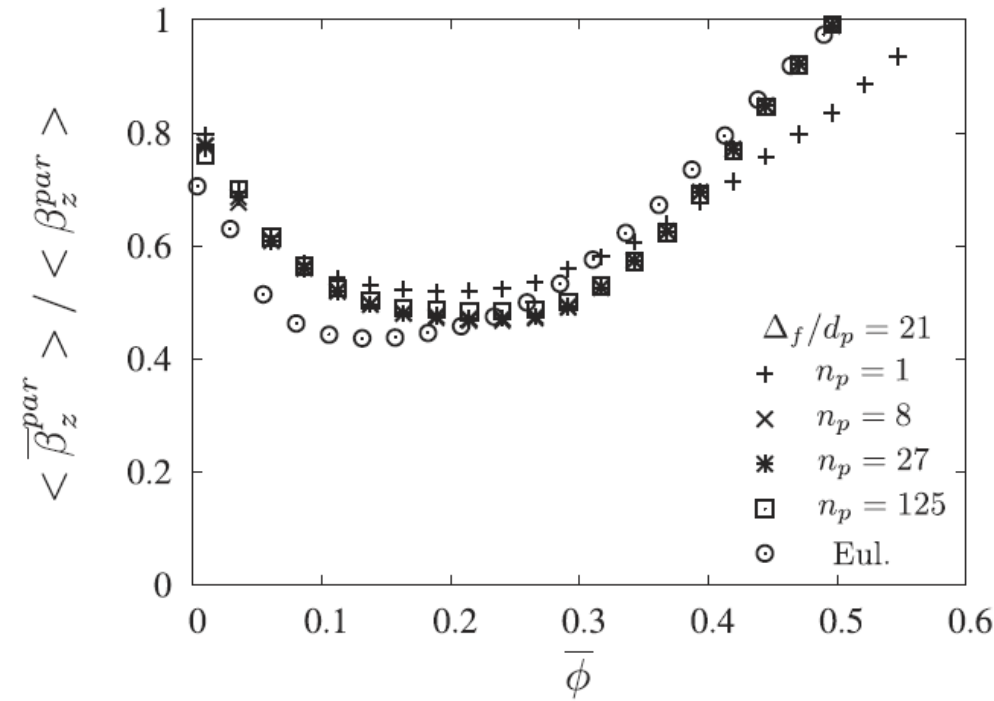


- CFD-DEM **allows two-step coarsening approach**
- In addition: use **projected (mean) particle speed**, which is approximating data from a Two-Fluid-Model (TFM).
- **3 choices of “filtered” coefficients!**

Q: What are the differences?

Faster Simulations

Deriving Closures for Large(r) Scale-Models



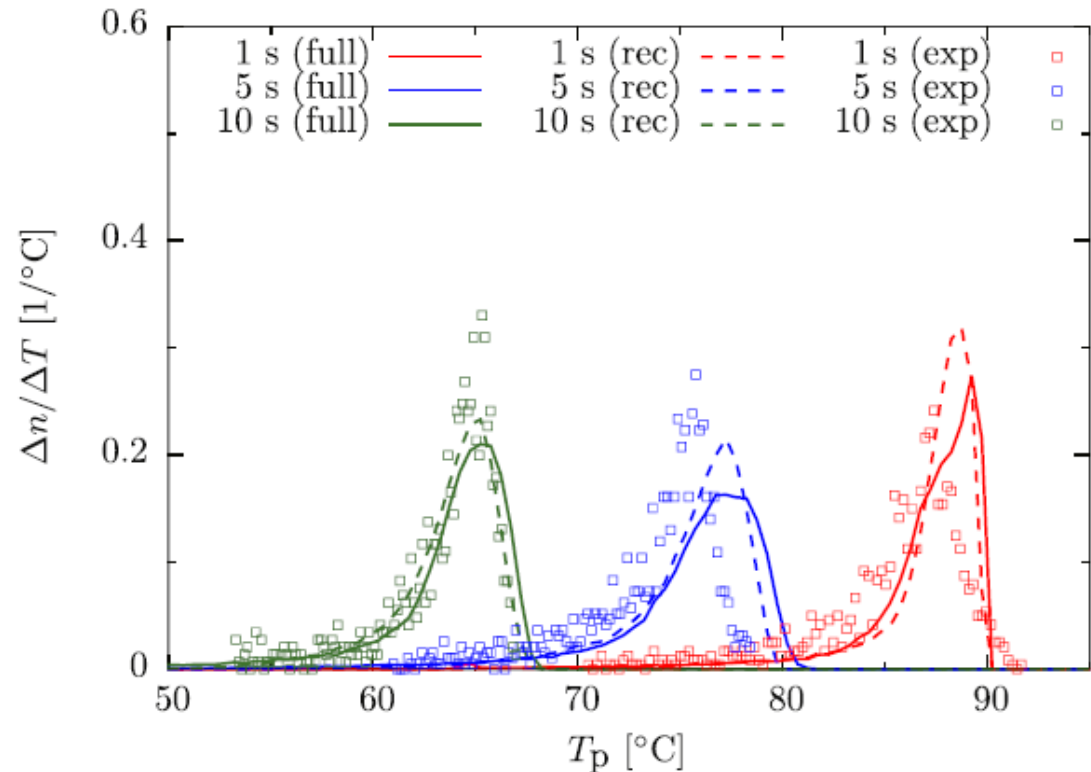
- **Fluid-coarsening** causes *the* dominant reduction in the effective drag coefficient
- **Only small difference** when going from coarse-grid CFD-DEM ($n_p = 1$) to parcel-based coarse-grid CFD-DEM
- Even **differences to TFM** are small!

A: only resolving the gradients in the voidage field is important (in line with theory [10])!

Even Faster Simulations

The long-standing Problem: How to put “Particles in the Loop”

- **Return to our starting point: ‘Meaningful** reaction kinetics must be fed into “micro-scale” models’
- **Option 1:** Direct approach: necessitates full simulation of all involved particles for long times (can be hours real-time!). **Not feasible for “closure screening”**
- **Option 2: “Record and Playback”** as suggested by Lichtenegger et al. [11] appears to be attractive! Currently demonstrated for **heat transfer (i.e., particle temperature distributions)**, but this could be extended.



Conclusions

- Closures for drag and heat/mass transfer are still **poor on a per-particle level** (and we even have not started looking at non-spherical or irregular particles!). **Particle (thermal) inertia “irons out” this problem**. But it persists for low particle-to-fluid density ratios, heat-sensitive reactions, etc.!
- **A first set of near wall corrections** ready to use! ...but there are still many improvements necessary near walls (e.g., wall-fluid heat transfer rates, polydispersity)
- A large number of closures need potential improvement. **Which one to attack first (experiments, DNS, calibration)?** Sensitivity analysis using fast meso-scale models appears essential.

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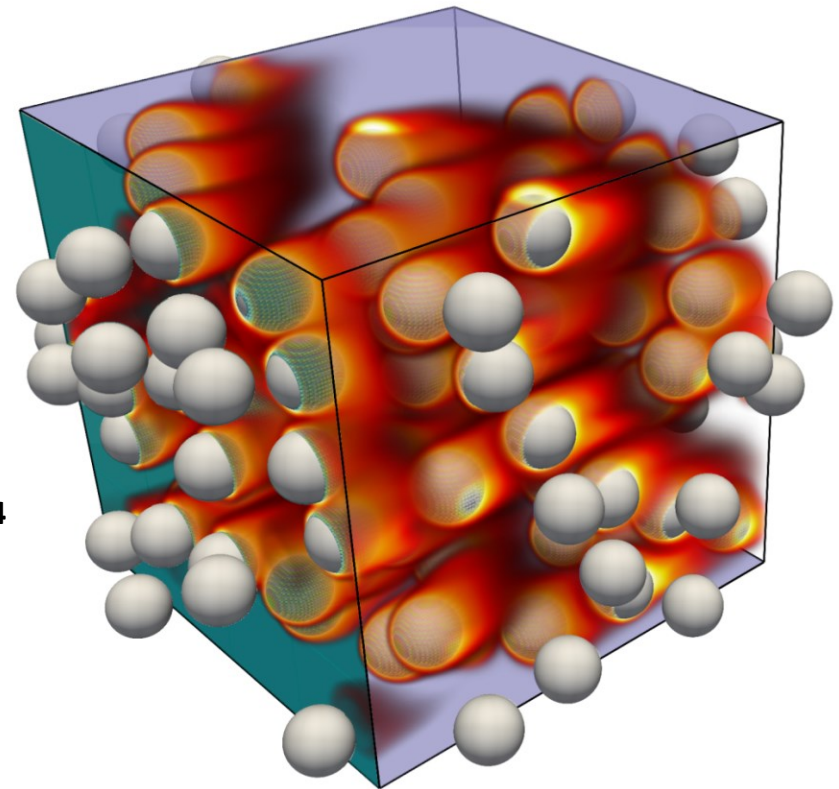
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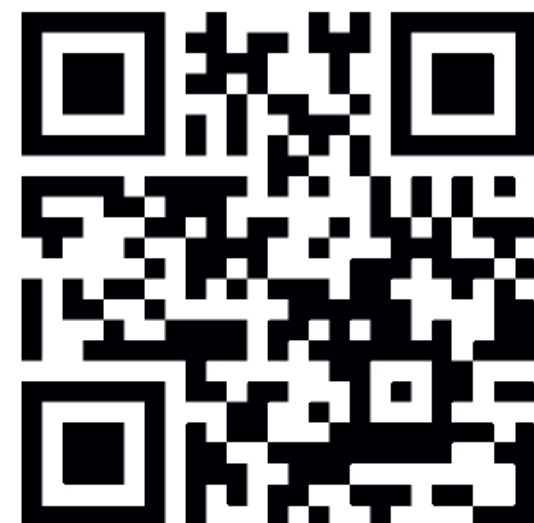


escape

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Parts of the “CPPPO” code were developed in the frame of the “NanoSim” project funded by the European Commission through FP7 Grant agreement no. 604656.

<http://www.sintef.no/projectweb/nanosim/>



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