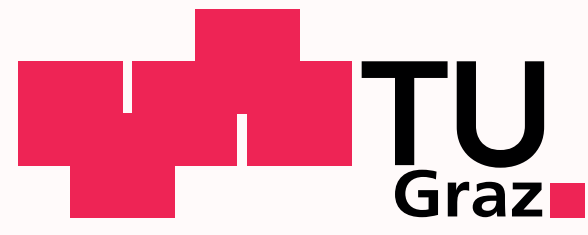


Variational Cluster Approach to the Bilayer Hubbard Model



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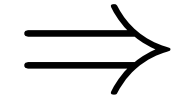


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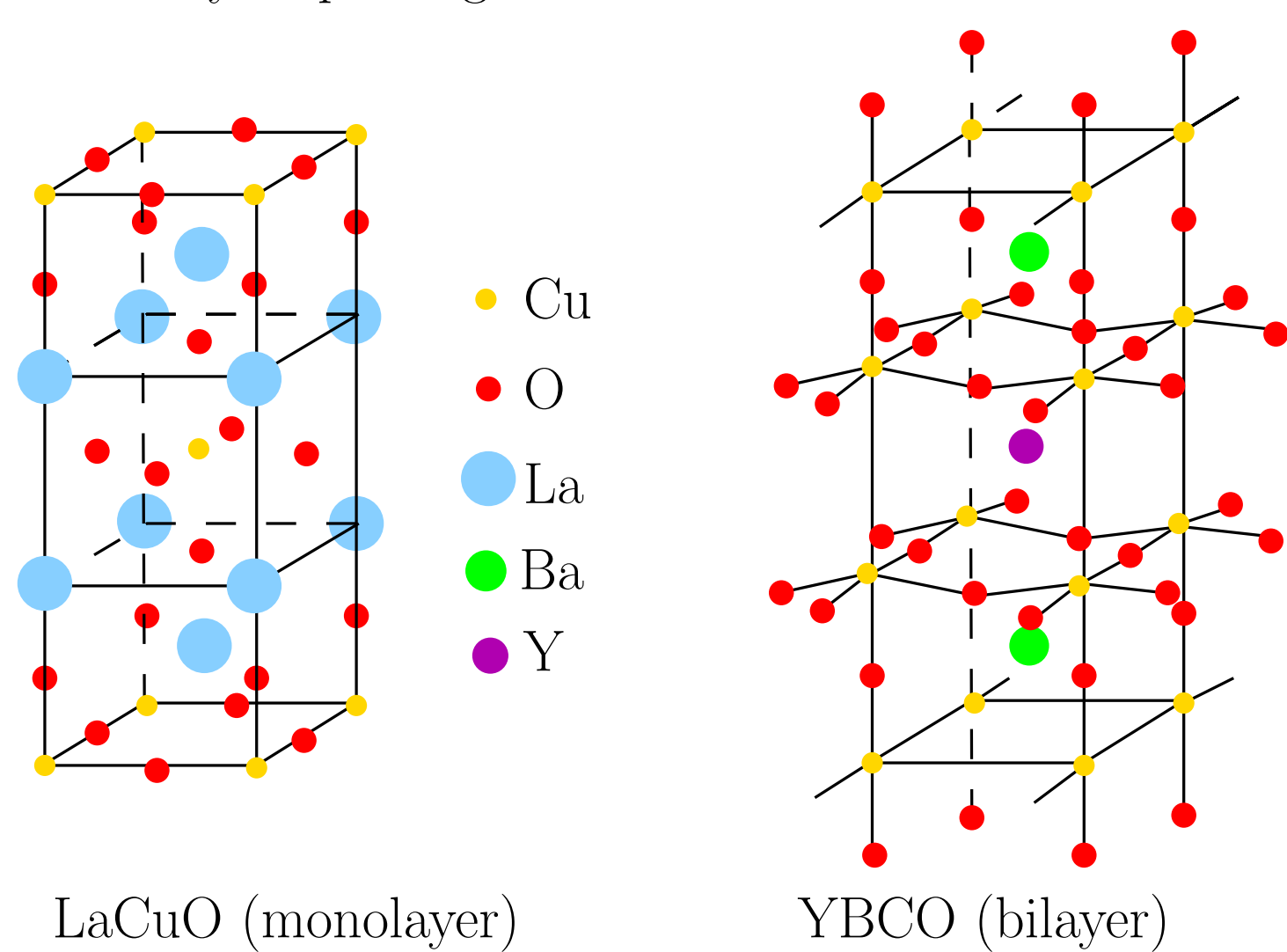
This work was supported by the Austrian Science Fund (FWF) under contract number P18551-N16

Bilayer High Temperature Superconducting Cuprates

- HTSC have perovskite structure
- superconductivity occurs in CuO_2 layers
- most HTSC contain two or more CuO_2 layers next to each other
- the interlayer hopping is expected to influence the critical temperature and doping dependence of superconductivity.
- ARPES experiments [1] show a bilayer splitting in the spectral functions of hole doped Bi2212, which is reduced from 88 meV in the normal state to 20 meV superconducting state.

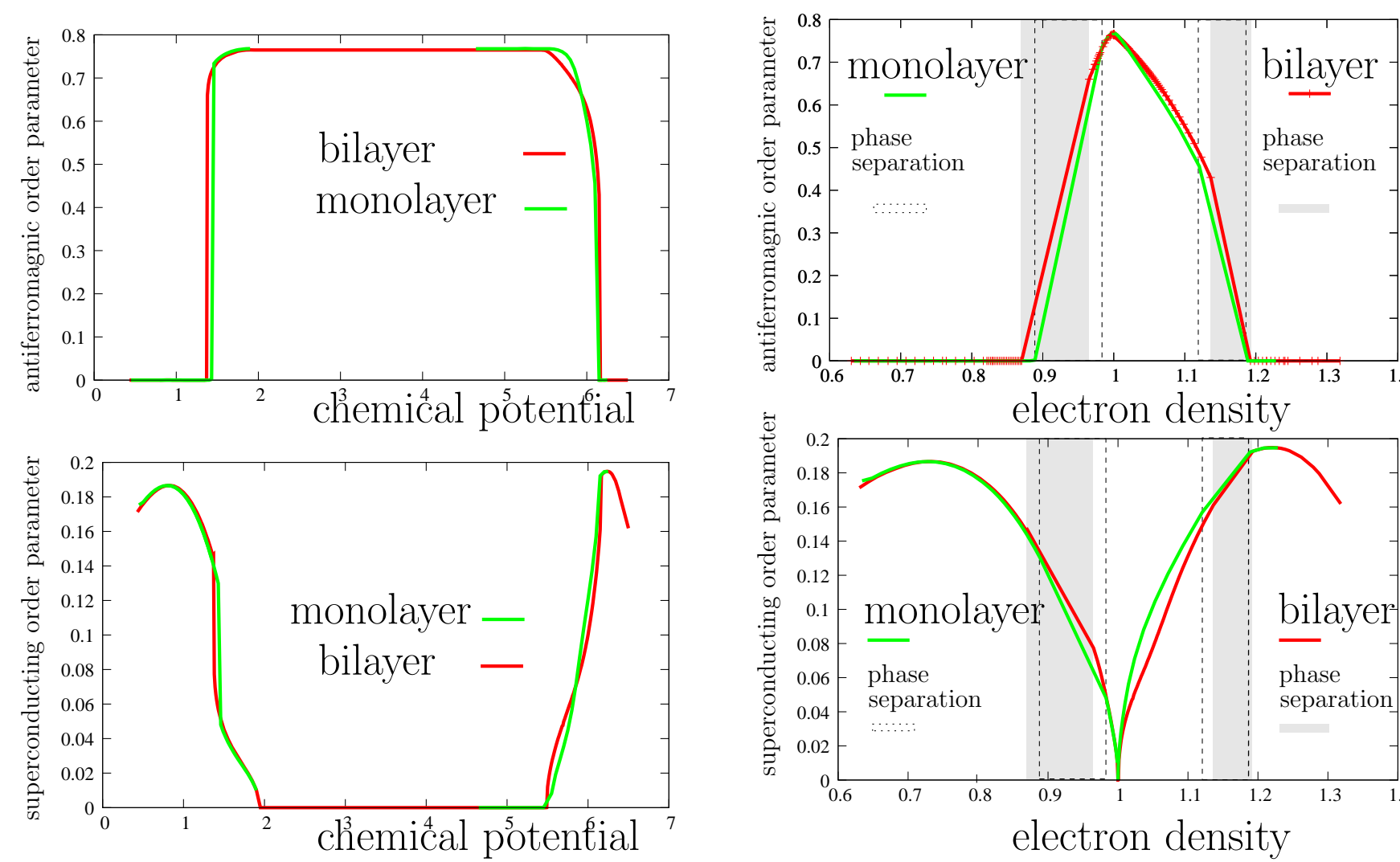


- use the Hubbard Model to describe a CuO_2 layer
- Variational Cluster Approach (VCA) ([2],[3]) to study phase diagram and spectral functions at different doping levels with and without interlayer hopping
- measure bilayer splitting



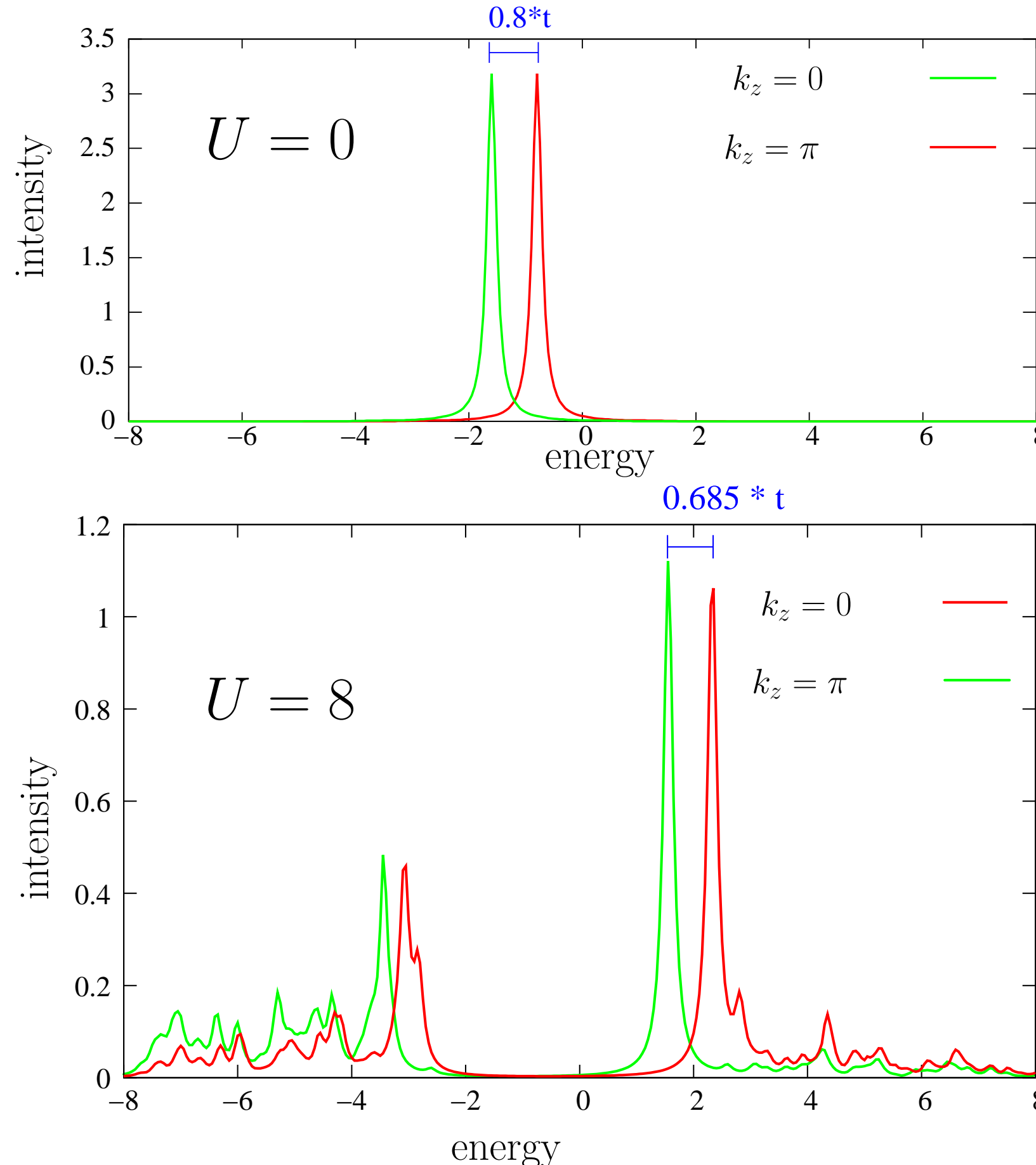
Results - Phase Diagram

- very similar to mono-layer phase diagram
- antiferromagnetic zone around half filling([8],[9],[10])
- at electron and hole doping a d-superconducting phase forms
- critical doping for superconductivity is different for electron and hole doping, as well as the vanishing of the antiferromagnetic phase.
- optimal doping is overestimated in size



Bilayer Splitting at Half Filling

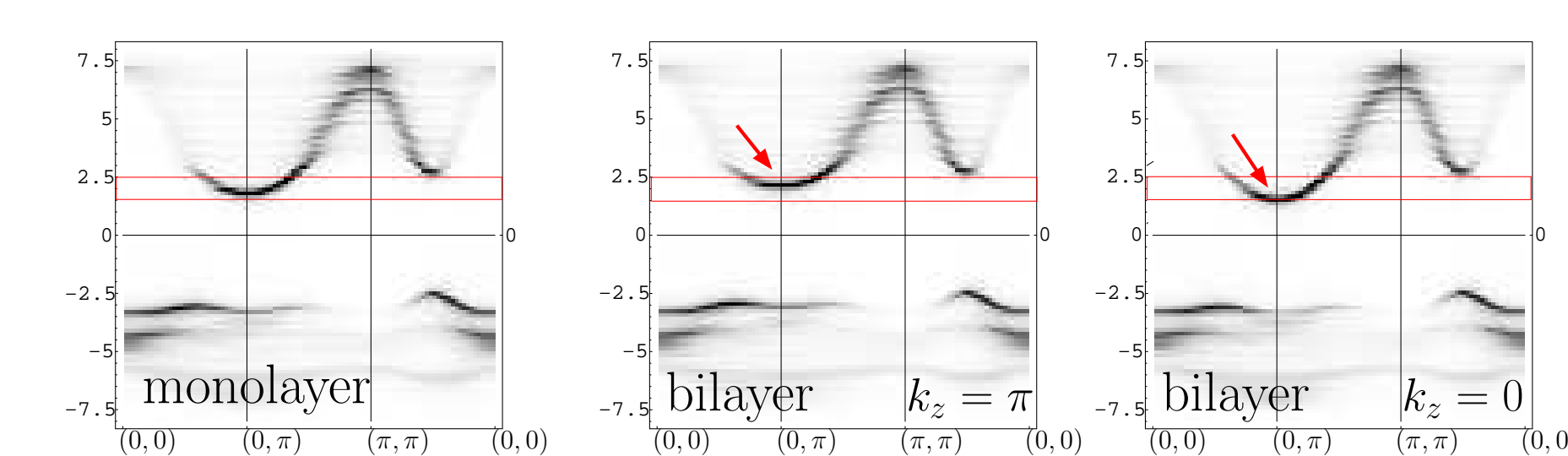
Density of states (DOS) at the point of maximum splitting $(0, \pi)$, for correlation $U = 0$ and $U = 8$.



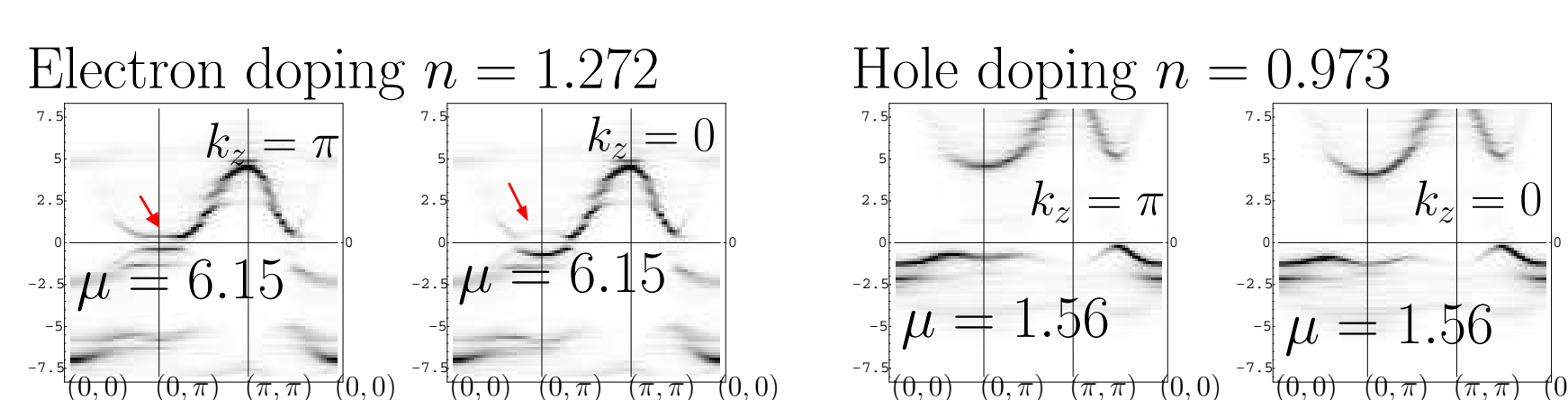
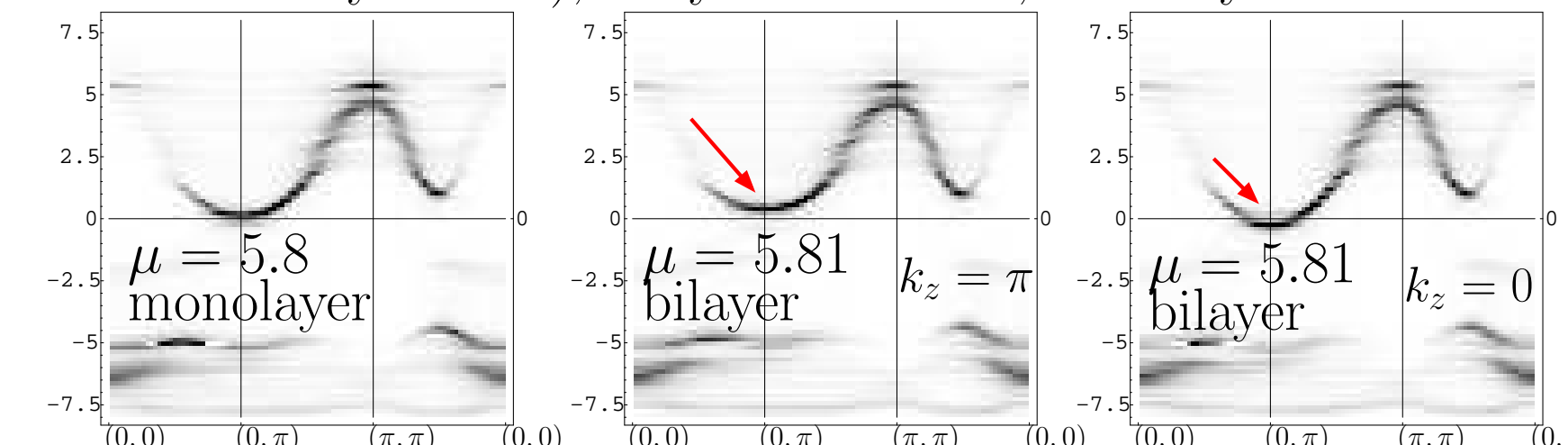
For $U = 0$ there is one excitation for $k_z = 0$ and $k_z = \pi$ each, with a splitting of $\Delta E(0, \pi) = 0.8t$.
For $U = 8$ the DOS is gapped and the splitting is reduced to approximately $\Delta E(0, \pi) = 0.685t$.

Results - Spectral Function

The spectral functions at half filling for monolayer and bilayer ($k_z = 0, k_z = \pi$) case:

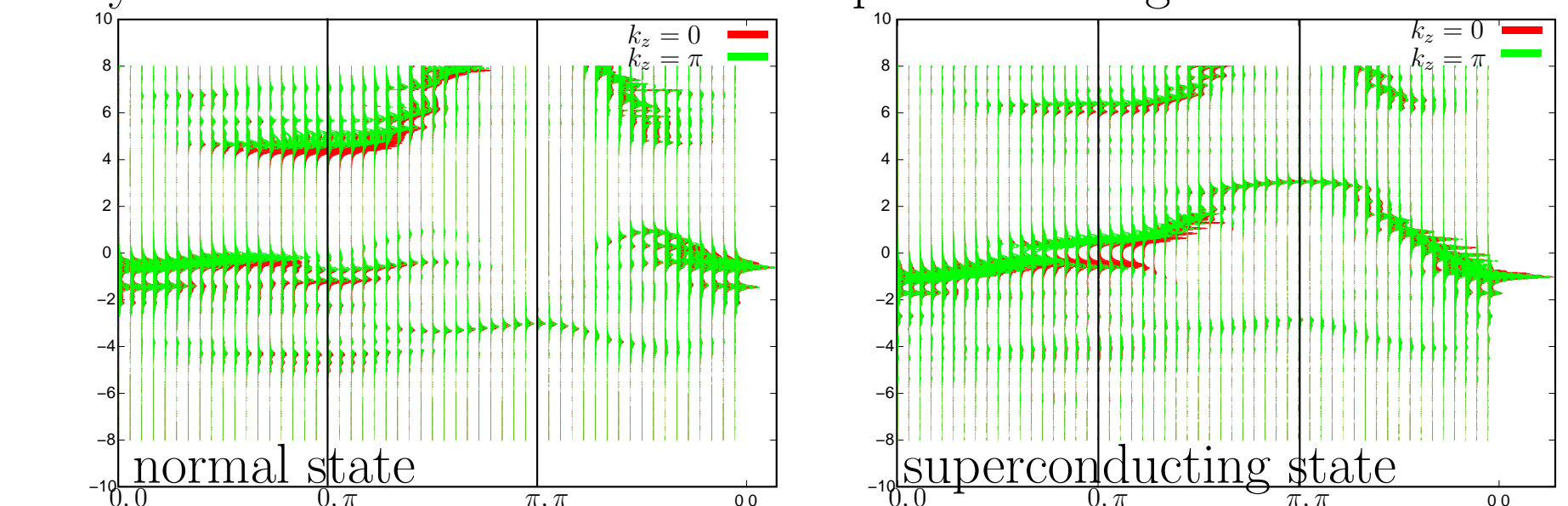


The spectral functions at electron doping (which in nature was not found for bilayer HTS), bilayer $n = 1.074$, monolayer $n = 1.03$:



Heavily Overdoped (With Holes)

- ARPES measurements on Bi2212 [1] indicate that the splitting is reduced in the superconducting case.
- \Rightarrow calculate the spectral function of the highly overdoped bilayer system both in the normal and superconducting state

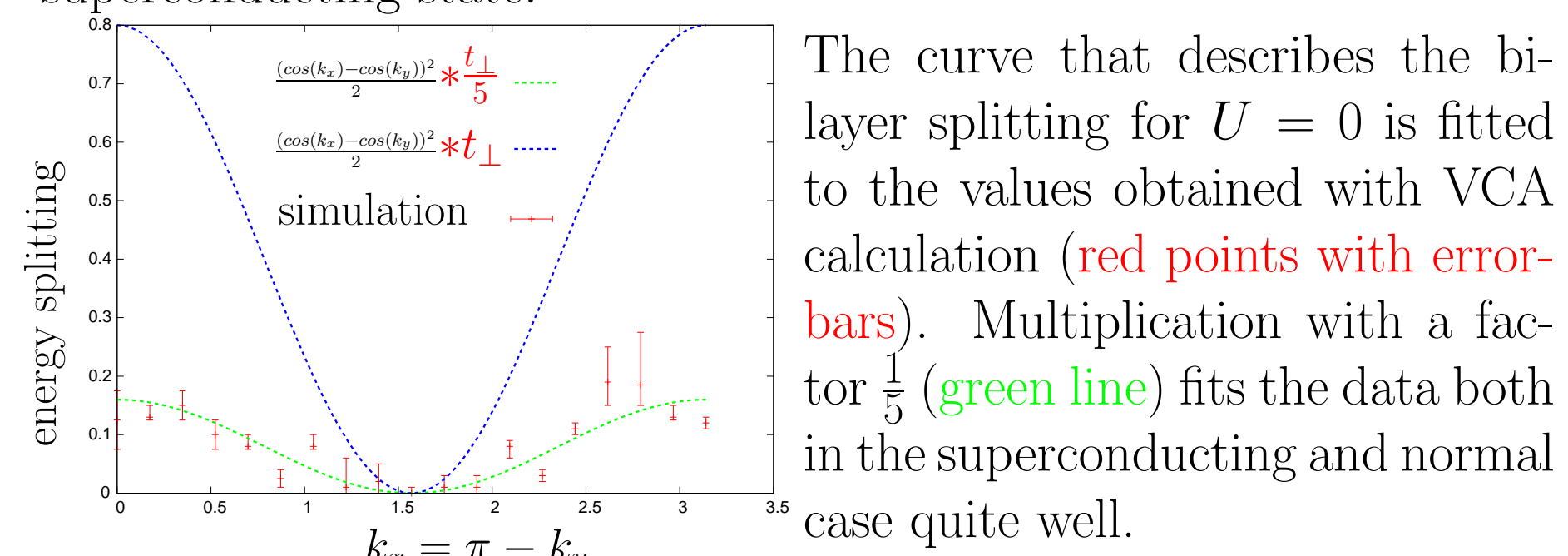


- compare bilayer splitting on the $\mathbf{k} = (0, \pi)$ to $\mathbf{k} = (\pi, 0)$ line of the Brillouin Zone for superconducting state and normal state with the values

$$\Delta E = \frac{(\cos(k_x) - \cos(k_y))^2}{2} t_{\perp} \quad (4)$$

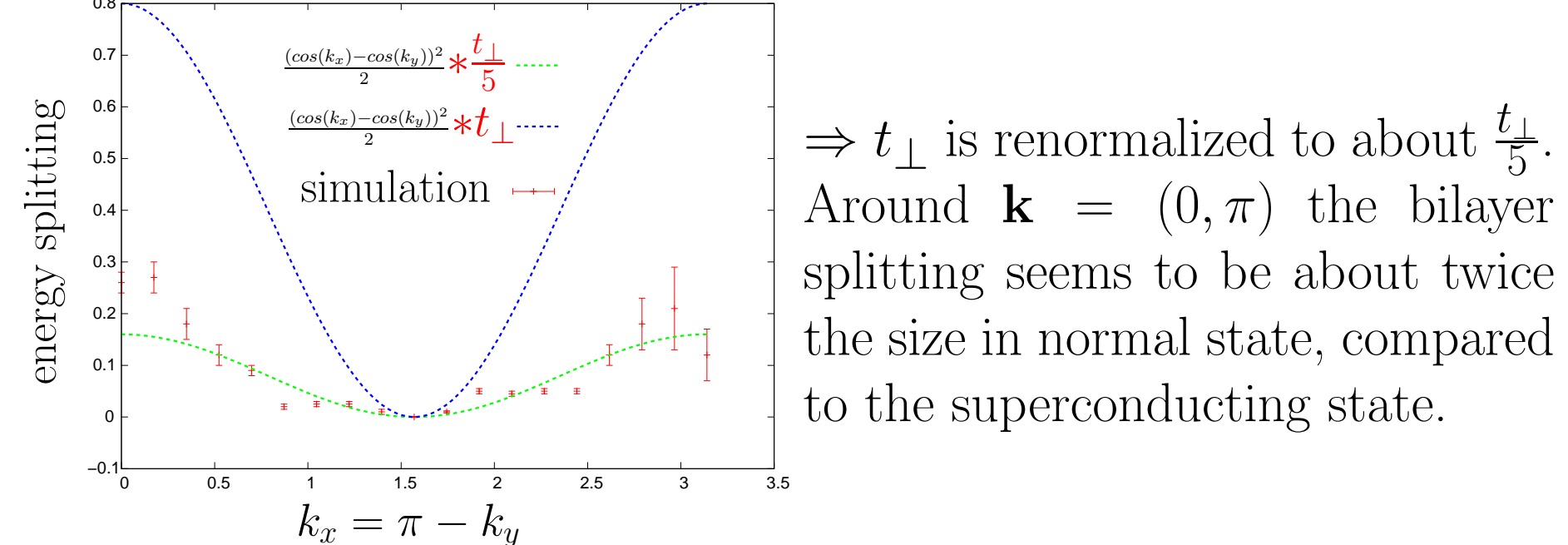
(blue line in figures), expected without correlation ($U = 0$).

superconducting state:



The curve that describes the bilayer splitting for $U = 0$ is fitted to the values obtained with VCA calculation (red points with error bars). Multiplication with a factor $\frac{1}{5}$ (green line) fits the data both in the superconducting and normal case quite well.

normal state:



$\Rightarrow t_{\perp}$ is renormalized to about $\frac{t_{\perp}}{5}$. Around $\mathbf{k} = (0, \pi)$ the bilayer splitting seems to be about twice the size in normal state, compared to the superconducting state.

Model

- Hubbard Hamiltonian (one layer)

$$H_H = -t \sum_{\langle ij \rangle, \sigma} (c_{i\sigma}^{\dagger} c_{j\sigma} + c_{j\sigma}^{\dagger} c_{i\sigma}) + t' \sum_{\langle\langle ij \rangle\rangle, \sigma} (c_{i\sigma}^{\dagger} c_{j\sigma} + c_{j\sigma}^{\dagger} c_{i\sigma}) + U \sum_i n_{i\uparrow} n_{i\downarrow} + \mu \sum_{i\sigma} n_{i\sigma} \quad (1)$$

- Interlayer hopping (in momentum space)[4]

$$H_{inter} = \sum_{k, \sigma} t_{\perp} \frac{v^2}{(1 - 2u t_{\perp}^2)} (c_{k\sigma a}^{\dagger} c_{k\sigma b} + c_{k\sigma b}^{\dagger} c_{k\sigma a}) \quad (2)$$

with $u = (\cos(k_x) + \cos(k_y))/2$ and $v = (\cos(k_x) - \cos(k_y))/2$, while a and b denote the layers. For $t'/t \ll 1$ it follows that

$$H_{inter} \approx \sum_{k, \sigma} t_{\perp} v^2 (c_{k\sigma a}^{\dagger} c_{k\sigma b} + c_{k\sigma b}^{\dagger} c_{k\sigma a}).$$

In real space this corresponds to

$$\sum_{\Delta x, \Delta y, \sigma} \frac{t_{\perp}}{4} (\delta_{\Delta x, 0} \delta_{\Delta y, 0} - \frac{1}{2} \delta_{\Delta x, \pm 1} \delta_{\Delta y, \pm 1} + \frac{1}{4} \delta_{\Delta x, 0, \pm 2} \delta_{\Delta y, 0, \pm 2}) (c_{i+\Delta x, \Delta y, \sigma}^{\dagger} c_{i\sigma} + h.c.).$$

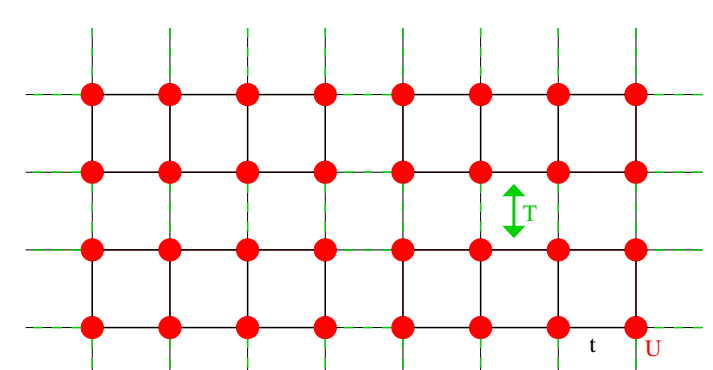
Parameters

U	$8t$
t'	$0.3t$
t_{\perp}	$-0.4t$

We have chosen $t_{\perp} = -0.4t$.

[4] get for YBCO $t_{\perp} \approx -0.625t$ and [5] use $t_{\perp} = -0.3t$ for BSC0.

Solver: Variational Cluster Approach



CPT:
 $H = H_{CL} + H_{intercl}$
 $G_{CPT}^{-1} = G_{CL}^{-1} - T$

- Exact diagonalization of the cluster Hamiltonian using the Lanczos procedure. (4x2x1 sites)
- Intercluster hoppings, including t_{\perp} , treated perturbatively (CPT,[6]).
- Variational calculation to allow symmetry-broken phases ([3],[2],[7]):

$$H'_{CL} = H_{CL} + h_{field} \quad H'_{intercl} = H_{intercl} - h_{field} \text{ with}$$

$$h_{field} = h_{SC} \sum_{i,j} \frac{n_{i,j}}{2} (c_{i,\uparrow} c_{j,\downarrow} + h.c.) + h_M \sum_{i\sigma} (-1)^{\sigma} e^{i\mathbf{Q}\cdot\mathbf{r}_i} c_{i,\sigma}^{\dagger} c_{i,\sigma} \quad (3)$$

Conclusions

The interlayer hopping

- introduces differences in the phase diagram between the monolayer and bilayer Hubbard model
- causes a splitting of the energy bands.

Comparison of the energy splitting for

$U = 0$

$U = 8$

- all bands split according to $t_{\perp} \frac{(\cos(k_x) - \cos(k_y))^2}{2}$
- Splitting of the bands has the same form but renormalized t_{\perp}
- heavily overdoped sample, superconducting state $\Delta E(0, \pi) \approx (0.15 \pm 0.05)t \approx 0.375t_{\perp}$
- heavily overdoped sample, normal state $\Delta E(0, \pi) \approx (0.26 \pm 0.04)t \approx 0.65t_{\perp}$
- maximum splitting amounts to $\Delta E(0, \pi) = 0.8t = 2t_{\perp}$
- half filling (antiferromagnetic state) $\Delta E(0, \pi) \approx 0.685t = 1.7125t_{\perp}$.

At $\mathbf{k} = (0, \pi)$ we see a decrease in bilayer splitting from normal to superconducting state of about 50%, which is much less than what obtained within experiment [1]

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